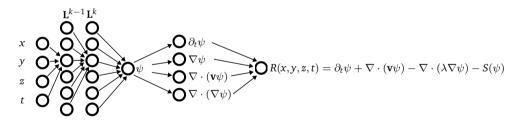
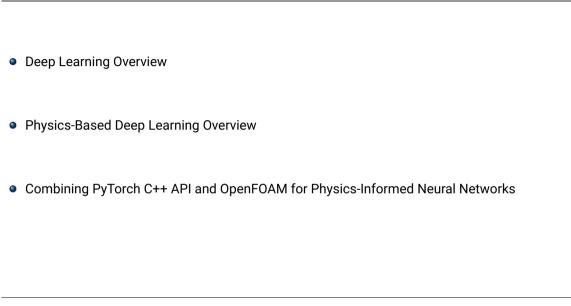
# Towards physics-based deep learning in OpenFOAM: Combining OpenFOAM with the PyTorch C++ API



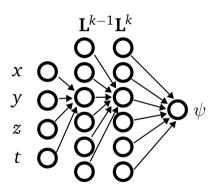
Tomislav Maric (TU Darmstadt), Andre Weiner (TU Braunschweig) 17th OpenFOAM Workshop, 11.07.2022, Cambridge University





# Deep Learning Neural Network - NN

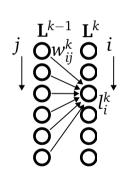




■ Neural Network (NN) has an input layer (e.g. (x,y,z,t)), D-1 hidden layers, and a (e.g. scalar) output layer  $\psi$ .

**Hidden Layers** 





Hidden layer  $\mathbf{L}^k$  is computed from the previous layer

$$l_i^k = \sum_{i=1}^{N_{L^{k-1}}} w_{ij}^k l_j^{k-1} + b_i^k \tag{1}$$

Einstein's notation (repeated index  $\equiv$  dot product)

$$l_i^k = w_{ij}^k l_j^{k-1} + b_i^k (2)$$

Matrix-vector product

$$\mathbf{L}^k = \mathbf{w}^k \cdot \mathbf{L}^{k-1} + \mathbf{b}^k \tag{3}$$

 $\mathbf{w}^k$  is a  $N_{\mathbf{L}^k} \times N_{\mathbf{L}^{k-1}}$  matrix.

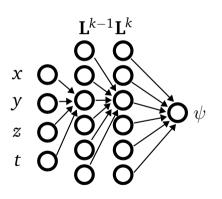


Adding activation functions to the layers results in the final NN, as a composition of functions

$$\psi_{\theta}^{nn}(\mathbf{u}) = \mathbf{w}^{D} \cdot L^{D-1} + b^{D} = \mathbf{w}^{D} \cdot \sigma(\mathbf{w}^{D-1} \cdot L^{D-2} + \mathbf{b}^{D-1}) + b^{D}$$

$$= \mathbf{w}^{D} \cdot \sigma(\mathbf{w}^{D-1} \cdot \sigma(\mathbf{w}^{D-2} \cdot \sigma(\dots \sigma(\mathbf{w}^{1} \cdot \mathbf{u} + \mathbf{b}^{1}) \dots) + \mathbf{b}^{D-2}) + \mathbf{b}^{D-1}) + b^{D}$$
(4)

- $\mathbf{u} = (x, y, z, t)$  in our example.
- m heta are all the weights and biases,  $heta=\{w_{ij}^k,b_i^k\}$ ,  $k\in[1,D]$ ,  $j\in[1,N_{\mathbf L^{k-1}}]$ ,  $i\in[1,N_{\mathbf L^k}]$ .
- When approximating functions, the last layer is often "linear".



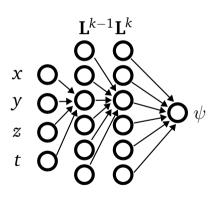
A set of points  $\{\mathbf{u}_p\}_{p\in P}$  and their data  $\{\psi_p\}_{p\in P}$  can be used to define an error of  $\psi_n^{nn}(\mathbf{u}_p)$ , e.g.

$$e_{MSE}(\theta) = \frac{1}{N_p} \sum_{p=1}^{N_p} (\psi_{\theta}^{nn}(\mathbf{u}_p) - \psi_p)^2.$$
 (5)

The network "learns" some  $\theta_M$  that minimize  $e_{MSE}$ .

$$\theta_M = \arg\min_{\theta} e_{MSE}(\theta) \tag{6}$$





•  $\theta_M = \arg\min_{\theta} e_{MSE}(\theta)$  requires  $\min_{\theta} e_{MSE}(\theta)$ 

$$N_p$$

 $\partial_{\theta} e_{MSE}(\theta) \rightarrow 0$ 

$$\sum_{p=1}^{N_p} (\psi_{\theta^m}^{nn}(\mathbf{u}_p) - \psi_p) \partial_{\theta_i^m} \psi_{\theta^m}^{nn}(\mathbf{u}_p) \stackrel{!}{=} 0$$
 (7)

- An approximation is inexact so generally  $\psi_{\theta^m}^{nn}(\mathbf{u}_p) \psi_p \neq 0$ , and we strive for  $\partial_{i_i}^m \psi_{\theta^m}^{nn}(\mathbf{u}_p) \to 0$ .
- Why m? In the beginning  $\theta^{m-1}$  is somehow (randomly) initialized and we iteratively (m) improve it to satisfy eq. (7).

Gradient descent



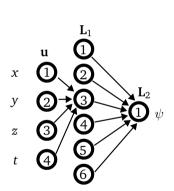
$$\sum_{p=1}^{N_p} (\psi_{ heta^m}^{nn}(\mathbf{u}_p) - \psi_p) \partial_{ heta_i^m} \psi_{ heta^m}^{nn}(\mathbf{u}_p) \stackrel{!}{=} 0$$

■ Imagine we somehow know  $\partial_{\theta_i^m} \psi_{\theta^m}^{nn}(\mathbf{u}_p)$  ( $\nabla_{\theta^m} \psi_{\theta^m}^{nn}(\mathbf{u}_p)$  in vector notation), we can then use gradient descent

$$\theta^{m+1}(\mathbf{u}_p) = \theta^m(\mathbf{u}_p) - \lambda^m \nabla_{\theta^m} \psi_{\theta^m}^{nn}(\mathbf{u}_p)$$
(8)

ensuring  $e_{\mathit{MSE}}(\theta^{m+1}) \leq e_{\mathit{MSE}}(\theta^m)$  (hopefully  $\|\nabla_{\theta^{m+1}}\psi^{nn}_{\theta^m}(\mathbf{u}_p)\|_2 \leq \|\nabla_{\theta^m}\psi^{nn}_{\theta^m}(\mathbf{u}_p)\|_2$ ), and adapting  $\lambda$  at m to tune the step size.

- New parameters are set as  $\frac{1}{N_p} \sum_{p=1}^{N_p} \theta^{m+1}(\mathbf{u}_p)$  (average), or by batch-average, or using radom-subsets of data points.
- Real-world algorithms adapt  $\lambda^m$  differently and are more complex.

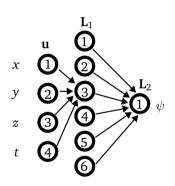


If 
$$\psi_{\theta}^{nn}(\mathbf{u}_p)=\mathbf{w}^2\cdot\sigma(\mathbf{w}^1\cdot\mathbf{u}+\mathbf{b}^1)+b^2$$
 what is  $\partial_{\mathbf{w}_{13}^1}\psi_{\theta}^{nn}(\mathbf{u}_p)$ ?

## **Deep Learning**

Gradient evaluation





lf

$$\psi_{\theta}^{nn}(\mathbf{u}_{p}) = \mathbf{w}^{2} \cdot \sigma(\mathbf{w}^{1} \cdot \mathbf{u} + \mathbf{b}^{1}) + b^{2}$$

what is  $\partial_{w_{12}^1} \psi_{\theta}^{nn}(\mathbf{u}_p)$ ?

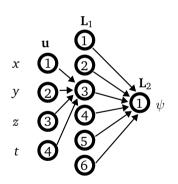
$$\partial_{w_{13}^1} \psi_{\theta}^{nn}(\mathbf{u}_p) = \mathbf{w}^2 \cdot \sigma'(\mathbf{w}^1 \cdot \mathbf{u} + \mathbf{b}^1) u_3$$

- Imagine writing this down for every  $w_{ij}^k, b_i^k$  for a deep NN.
- To make things worse, the number of layers and nodes change during "hyperparameter tuning".

## **Deep Learning**

**Gradient evaluation - Finite Differences (FD) I** 





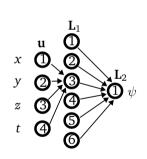
Finite Differences (FD) generalize to arbitrary NN architectures, but don't work, because of **computational costs**.

- $\mathbf{w}^k$  is a  $N_{\mathbf{L}^k} \times N_{\mathbf{L}^{k-1}}$  matrix.
- **b**<sup>k</sup> is a  $N_{\mathbf{L}^k}$  vector.
- lacksquare For  $\partial_{ heta_i^m} \psi_{ heta^m}^{nn}(\mathbf{u}_p)$ , we need

$$N_{ heta} := \sum_{i=2}^{D} N_{\mathbf{L}^{k}} (1 + N_{\mathbf{L}^{k-1}})$$
 (9)

finite differences, one for each weight and bias.





Finite Differences (FD) generalize to arbitrary NN architectures, but don't work, because of **Floating-Point (FP)** cancellation errors.

- As we converge towards  $\theta^M$ , forward passes get **very close to** each other  $|\psi_{\theta^m}^{nn}(\mathbf{u}_p) \psi_{\theta^{m-1}}^{nn}(\mathbf{u}_p)| \to 0$ .
- The computer has limited precision, so as weights  $\theta_i^m, \theta_i^{m-1}$  get close to each other

```
theta_i_m = 1.1234567891234569|33333333333
theta_i_m_1 = 1.1234567891234102|2222222222
```

their difference



#### Some Finite Differences

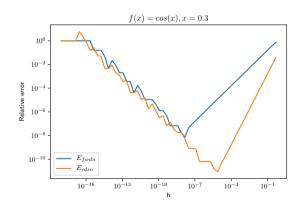
$$f'_{fwd} = \frac{f(x+h) - f(x)}{h} + O(h) \tag{10}$$

$$f'_{cds} = \frac{f(x+h) - f(x-h)}{2h} + O(h^2)$$
 (11)

**Important:** Finite Differences are **inexact**, order-of accuracy  $O(h^p)$ . A great book on FPA is Overton [2001]

- Once  $h \le 0.5ulp(x)$  with nearest rounding, full cancellation occurs.
- ulp units in the last place,  $ulp(x) = 2^{-52}2^E$ , E is the exponent.





### Relative derivative (gradient) error

$$e_n^{fwd,cds} = \frac{|f'_{fwd,cds} - f'_{exact}|}{|f'_{exact}|}.$$
 (12)

 $h \to 0$ ,  $e_n^{fwd,cds} \to 1$ : Floating-Point cancellation errors prevent asymptotic convergence of Finite Differences.

## **Deep Learning**

**Gradient evaluation - Symbolic Calculation** 

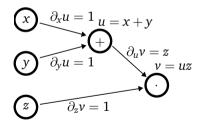


- Instead of manually evaluating  $\partial_{\theta_i^m} \psi_{\theta^m}^{nn}(\mathbf{u}_p)$ , let some software (like Sympy) write down closed-form expressions for you using Symbolic Calculation.
- Doesn't work: huge closed-form expressions are necessary, causing huge memory and CPU overheads.



$$f(x,y,z) = (x+y)z = f(v(u(x,y),z))$$

Notation: 
$$\partial_s f = \frac{\partial f(s)}{\partial s}$$



Reverse-mode Automatic Differentiation (AD, details in Griewank and Walther [2008]) is the basis for evaluating derivatives for NN training (backpropagation).

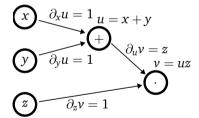
- Mathematic expressions are modeled with an directed acyclic graph (DAG).
- Intermediate results stored in variables.
- The graph's edges can evaluate known partial derivatives w.r.t. intermdediate variables.
- Chain rule is used to compute the partial derivative along the graph:

$$\partial_x f = \partial_v f(v) \partial_u v(u, z) \partial_x u(x, y) = 1z1 = z$$
 (13)



$$f(x,y,z) = (x+y)z = f(v(u(x,y),z))$$

Notation: 
$$\partial_s f = \frac{\partial f(s)}{\partial s}$$



- **Exact**: no Finite Difference-induced Floating-Point cancellation errors, no discretization errors.
- Automatic for arbitrary NN architecture.
- Computationally more efficient than Symbolic Calculations or Finite Differences.
- Responsible for "reviving" Deep Learning.



- Math proves  $\theta^M$  exists and NNs are universal function approximators, but not how to find it.
- Finding  $\theta^M$  depends on  $\lambda$ , the NN architecture, and the activation function **hyperparameters**.
- Hyperparameters are "free" parameters tuned by
  - graduate/Ph.D. students (student descent algorithm)
  - Grid Search, Monte Carlo, Bayesian Optimization: keyword AutoML.
- Once hyperparameters are "tuned", **some**  $\theta^M$  is found with a minimal  $e_{MSE}$  in the best case over a response-surface that **hopefully** models the hyperparameter space well there is no guarantee  $\theta^M$  is globally-optimal in terms of data or hyperparameters.
- Training takes a lot of computational time and resources.
- As soon as a form of Stochastic Gradient Descent is used (large data), running the training twice with the same hyperparameters and input data will give a different output from the NN.

Summarv

- NNs are function compositions, composing a matrix/vector product, an addition, and a nonlinear (activation) function  $\sigma$ ,  $\psi_{\theta}^{nn}(\mathbf{u}) = \mathbf{w}^D \cdot \sigma(\mathbf{w}^{D-1} \cdot \sigma(\mathbf{w}^{D-2} \cdot \sigma(\dots \sigma(\mathbf{w}^1 \cdot \mathbf{u} + \mathbf{b}^1) \dots) + \mathbf{b}^{D-2}) + \mathbf{b}^{D-1}) + b^D$ , for our example scalar function  $\psi$ .
- An NN is a function-approximator, "trained" by minimizing an error norm over data, like MSE  $\sum_{p=1}^{N_p} (\psi_{\theta^m}^{nn}(\mathbf{u}_p) \psi_p) \partial_{\theta_i^m} \psi_{\theta^m}^{nn}(\mathbf{u}_p) \stackrel{!}{=} 0$
- Approximation generally means  $|\psi_{\theta^m}^{nn}(\mathbf{u}_p) \psi_p| \neq 0$ , we aim for  $\partial_{\theta_i^m}\psi_{\theta^m}^{nn}(\mathbf{u}_p) \stackrel{!}{=} 0$ .
- To reach this, we perform (some form of) gradient descent  $\theta^{m+1}(\mathbf{u}_p) = \theta^m(\mathbf{u}_p) \lambda^m \nabla_{\theta^m} \psi_{\theta^m}^{nn}(\mathbf{u}_p)$ , resulting in  $e_{MSE}(\theta^{m+1}) \leq e_{MSE}(\theta^m)$ .
- For gradient descent, we compute  $\partial_{\theta_i^m} \psi_{\theta^m}^{nn}(\mathbf{u}_p)$  (gradient components), using Reverse-mode Automatic Differentiation.

## **Physics-Based Deep Learning**

Literature survey (incomplete)



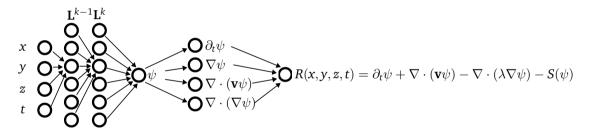
Different approaches exist, all extend the idea of function-approximation by NNs with satisfying PDEs. PDEs are built from differential operators, that are constructed from the NN using AD.

- The idea originated (afaik) with Lagaris et al. [1998].
  - A collocation method with NN as a trial function.
- Geometrically complex boundaries: Lagaris et al. [2000], McFall and Mahan [2009].
- Galerkin method with NN instead of shape functions: Sirignano and Spiliopoulos [2018].
- Raissi et al. [2019], Physics-Informed Neural Networks (PiNN)s collocation MSE for PDEs like Lagaris et al. [1998] + data MSE.
- More alternatives described by Thuerey et al. [2022].

#### This talk addresses PiNNs.

PiNNs one slide





- 1. Re-use Automatic Differentiation used for NN training, for computing partial derivatives of the forward-pass with respect to NN input to construct PDE operators.
- 2. Extend the loss function with PDE residuals: the NN learns data and the PDE.

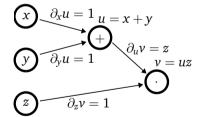
## **Physics-Based Deep Learning**

Re-using Automatic Differentiation for differential operators in PDEs



$$f(x,y,z) = (x+y)z = f(v(u(x,y),z))$$

Notation: 
$$\partial_s f = \frac{\partial f(s)}{\partial s}$$



- Evaluating  $\partial_{\theta_i^m} \psi_{\theta_i^m}^{nm}(\mathbf{u}_p)$  generates partial derivaties w.r.t. intermediate variables.
- The cached partial derivatives are re-used to compute e.g.  $\partial_t \psi_{\theta^m}^{nn}(\mathbf{u}_p)$ ,  $\mathbf{u}_p = (x,y,z,t)$ , and equivalently for x,y,z,

$$\nabla \psi_{\theta^m}^{nn}(\mathbf{u}_p) = \begin{bmatrix} \partial_x \psi_{\theta^m}^{nn}(\mathbf{u}_p) \\ \partial_y \psi_{\theta^m}^{nn}(\mathbf{u}_p) \\ \partial_z \psi_{\theta^m}^{nn}(\mathbf{u}_p) \end{bmatrix}$$
(14)

and have to be cached again for higher-order differential operators like  $\nabla \cdot \nabla \psi_{\theta^m}^{nn}(\mathbf{u}_p)$ .

Constructing the PiNN PDE residual I



- The residual of the PDE is computed using differential operators computed exactly<sup>1</sup> by the Automatic Differentiation framework there are no discretization errors.
- The residual of a passive scalar transport equation is

$$R_{\theta}(\mathbf{u}_p) = R_{\theta}(x, y, z, t) = \partial_t \psi_{\theta^m}^{nn}(\mathbf{u}_p) + \nabla \cdot (\mathbf{v}\psi_{\theta^m}^{nn}(\mathbf{u}_p)) - \nabla \cdot (\lambda \nabla \psi_{\theta^m}^{nn}(\mathbf{u}_p)) - S(\psi_{\theta^m}^{nn}(\mathbf{u}_p))$$
 (15)

evaluated at collocation points  $\{\mathbf{u}_p\}_{p\in[1,...,N_p}$ .

• Contrary to training the NN only on data with  $\partial_{\theta_i^m} \psi_{\theta^m}^{nn}(\mathbf{u}_p)$ , the PiNN residual requires partial derivatives of the NN forward-pass with respect to NN inputs.

<sup>&</sup>lt;sup>1</sup>Up to floating-point arithmetic errors.

Constructing the PiNN PDE residual II

- One technical detail on how the partial derivatives are evaluated by AD frameworks (e.g. torch::autograd) is relevant for constructing PiNNs in practice.
- The NN  $\psi : \mathbb{R}^k \to \mathbb{R}^l$ ,  $\psi = \psi(\mathbf{u}, \theta)$ ,  $k = d + 1 + N_\theta$ : d for spatial dimensions and +1 for time in  $\mathbf{u} = (x, y, z, t)$ ,  $N_\theta$  for all the weights and biases.
- lacksquare The loss  $e_{MSE}:\mathbb{R}^l
  ightarrow\mathbb{R}$ ,  $e_{MSE}=e_{MSE}(\psi(\mathbf{u}, heta))$ .
- The  $\partial_{\theta_i}e_{MSE}(\psi(\mathbf{u},\theta))$  for the gradient descent, is by chain rule in Einstein's notation

$$\frac{\partial \psi_i}{\partial \theta_i} \frac{\partial e_{MSE}}{\partial \psi_i} = \frac{\partial e_{MSE}}{\partial \theta_i},\tag{16}$$

in matrix notation

$$J_{\theta}\psi \cdot J_{\psi}e_{MSE}$$
 (17)

is a dot product of two Jacobians - the NN w.r.t  $\theta$ , the loss w.r.t. NN output.

Constructing the PiNN PDE residual III



For our example NN  $\psi: \mathbb{R}^4 \to \mathbb{R}$ ,  $\psi(\mathbf{u}) = \psi(x,y,z,t)$  (dropping  $^{nn}$  superscript),

$$J_{u}\psi(\mathbf{u}_{p}) = \begin{bmatrix} \partial_{x}\psi(\mathbf{u}_{p}) \\ \partial_{y}\psi(\mathbf{u}_{p}) \\ \partial_{z}\psi(\mathbf{u}_{p}) \\ \partial_{t}\psi(\mathbf{u}_{p}) \end{bmatrix} 1, \tag{18}$$

where 1 is the 0-rank tensor (scalar) we use to obtain the elements of  $\nabla \psi$  and  $\partial_t \psi$ ,

$$\nabla \psi = \begin{bmatrix} (J_u \psi(\mathbf{u}_p))_1 \\ (J_u \psi(\mathbf{u}_p))_2 \\ (J_u \psi(\mathbf{u}_p))_3 \end{bmatrix}, \quad \partial_t \psi = (J_u \psi(\mathbf{u}_p))_4$$
(19)

Constructing the PiNN PDE residual IV



Hessian matrix

$$(H_{u}\psi)^{T} = (J_{u}J_{u}\psi)^{T} = \begin{pmatrix} J_{u} \begin{bmatrix} \partial_{x}\psi \\ \partial_{y}\psi \\ \partial_{z}\psi \\ \partial_{t}\psi \end{bmatrix} \end{pmatrix}^{T} = \begin{bmatrix} \partial_{x}^{2}\psi & \partial_{x}\partial_{y}\psi & \partial_{x}\partial_{z}\psi\partial_{x} & \partial_{x}\partial_{t}\psi \\ \partial_{y}\partial_{x}\psi & \partial_{y}^{2}\psi & \partial_{y}\partial_{z}\psi & \partial_{y}\partial_{t}\psi \\ \partial_{z}\partial_{x}\psi & \partial_{z}\partial_{y}\psi & \partial_{z}^{2}\psi & \partial_{z}\partial_{t}\psi \\ \partial_{t}\partial_{x}\psi & \partial_{t}\partial_{y}\psi & \partial_{t}\partial_{z}\psi & \partial_{t}^{2}\psi \end{bmatrix}$$
(20)

so

$$\nabla \cdot \nabla \psi = \Delta \psi = H^T \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} [1110] = \partial_x^2 \psi + \partial_y^2 \psi + \partial_z^2 \psi \tag{21}$$

Is eq. (21) valid if  $\mathbf{u}_{p} = (t, x, y, z)$ ?

## **Physics-Based Deep Learning**

Constructing the PiNN PDE residual IV



Hessian matrix

$$(H_{u}\psi)^{T} = (J_{u}J_{u}\psi)^{T} = \begin{pmatrix} J_{u} \begin{bmatrix} \partial_{x}\psi \\ \partial_{y}\psi \\ \partial_{z}\psi \\ \partial_{t}\psi \end{bmatrix}^{T} = \begin{bmatrix} \partial_{x}^{2}\psi & \partial_{x}\partial_{y}\psi & \partial_{x}\partial_{z}\psi\partial_{x} & \partial_{x}\partial_{t}\psi \\ \partial_{y}\partial_{x}\psi & \partial_{y}^{2}\psi & \partial_{y}\partial_{z}\psi & \partial_{y}\partial_{t}\psi \\ \partial_{z}\partial_{x}\psi & \partial_{z}\partial_{y}\psi & \partial_{z}^{2}\psi & \partial_{z}\partial_{t}\psi \\ \partial_{t}\partial_{x}\psi & \partial_{t}\partial_{y}\psi & \partial_{t}\partial_{z}\psi & \partial_{t}^{2}\psi \end{bmatrix}$$
(20)

so

$$\nabla \cdot \nabla \psi = \Delta \psi = H^T \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} [1110] = \partial_x^2 \psi + \partial_y^2 \psi + \partial_z^2 \psi \tag{21}$$

Is eq. (21) valid if  $\mathbf{u}_p=(t,x,y,z)$ ? No - order of partial derivatives changes changes  $J_u\psi$  and  $H_u\psi$ .

## **Physics-Based Deep Learning**

PiNN activation functions



- Since we're differentiating  $\psi_{\theta}^{nn}(\mathbf{u_p})$  w.r.t  $\mathbf{u}_p$ , the activation functions should be differentiable.
- Differentiable activation functions increase the number training iterations.

**PiNN loss function** 

■ Total PiNN loss is a weighted sum of the data-loss  $e^d_{MSE}$  (at data points or collocation points), residual-loss  $e^r_{MSE}$ , boundary-condition loss  $e^b_{MSE}$ , and initial-value loss  $e^i_{mse}$ 

$$l_{MSE} = e_{MSE}^d + e_{MSE}^r + e_{MSE}^{bd} + e_{MSE}^{bn} + e_{mse}^i$$
 (22)

$$e_{MSE}^{d} = \frac{1}{N_{d}} \sum_{p=1}^{N_{d}} (\psi^{nn}(\mathbf{u}_{p}) - \psi_{p})^{2}, \quad e_{MSE}^{bd} = \frac{1}{N_{bd}} \sum_{q=1}^{N_{bd}} (\psi^{nn}(\mathbf{u}_{q}) - \psi_{q})^{2}$$

$$e_{MSE}^{bn} = \frac{1}{N_{bd}} \sum_{r=1}^{N_{bn}} (\nabla \psi^{nn}(\mathbf{u}_{r}))^{2}, \quad e_{MSE}^{i} = \frac{1}{N_{bd}} \sum_{r=1}^{N_{i}} (\psi^{nn}(x, y, z, t = t_{0}) - \psi_{p}(x, y, z, t = t_{0}))^{2}$$

$$(23)$$

• Minimizing the total loss  $l_{\it MSE}$  makes the PiNN satisfy data, the PDE, and initial/boundary conditions in the least-squares sense.



#### Inverse problems

**PiNN advantages** 

- Recover the solution of a PDE with partially-known boundary and initial conditions, and some noisy measurements.
- Learn a heat transfer coefficient from measurements.

#### **High-dimensional PDEs**

Solving high-dimensional PDEs (finance?, physics) - Finite Differences do not scale for this.

#### Optimization

 Once trained, a PiNN surrogate model is very fast to evaluate, and should be (way) more accurate than an 1D ROM ODE solution. ■ The PiNN loss  $l_{MSE} = e^d_{MSE} + e^r_{MSE} + e^{bd}_{MSE} + e^{bn}_{MSE} + e^i_{mse}$  errors have different values and often (very) different gradients.

$$l_{MSE} = \lambda_d e_{MSE}^d + \lambda_r e_{MSE}^r + \lambda_{bd} e_{MSE}^{bd} + \lambda_{bn} e_{MSE}^{bn} + \lambda_i e_{mse}^i$$
(24)

- yaay, more free (PiNN loss regularization) "free" parameters to guesstimate. Alternatives: adaptive activation functions by Jagtap et al. [2020], modifying Adam solver by Wang et al. [2021], meta-learning by Psaros et al. [2022].
- Different error contribution may mean e.g. that a PiNN with data-assimilation satisfies the data and the PDE residual better than the BCs and ICs.



- PiNNs inherit the curse of hyperparameter dimensionality from DL.
- PiNNs often take many iterations (epochs) to train: differentiable activation functions have diminishing gradients and require small optimizer step sizes.
- Evaluating many PDE differential operators increases computational costs per epoch.
- Yes, a forward pass of a trained NN generates PDE solutions very quickly, but it still takes a lot of (unautomated) effort to get there.

# **Physics-Based Deep Learning PiNN Summary**



PiNNs are a relatively simple tool worth looking at

- Understand what the NN approximates.
- Understand how the Jacobian is used to construct PDE differential operators.
- Write down the total loss as the sum of: (data), residual, internal, and boundary condition loss.
- Apply adaptive activation functions or something similar to improve training.

If you just want to solve a "normal" PDE, it is overall faster to use a classical numerical method - there is potential in combining Physics-Based ML with classical numerical methods.

# Physics-Based Deep Learning Other videos

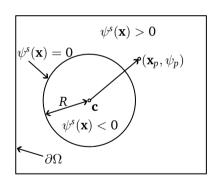


- "Introduction to Scientific Machine Learning 2: Physics-Informed Neural Networks", C. Rackauckas
- Hands-on introduction to Physics-Informed Machine Learning, I. Bilionis, A. Hans
- "When and why Physics-informed Neural Networks fail to train", P. Perdikaris
- "Physics-Informed Neural Networks | Misconceptions", C. Rackauckas

## A PiNN with OpenFOAM and PyTorch C++ API

TECHNISCHE UNIVERSITÄT DARMSTADT

Problem definition I



An implicit sphere is given as a zero-level set

$$\Sigma = \{ \mathbf{x} : \psi^e(\mathbf{x}) = 0 \}$$
 (25)

of a signed-distance field to a sphere

$$\psi^e(\mathbf{x}) = \|\mathbf{x} - \mathbf{c}\|_2 - R. \tag{26}$$

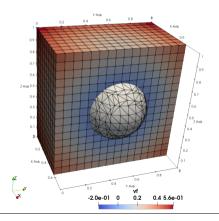
■ Approximate the data from points  $(\mathbf{x}_p, \psi_p^e)$  randomly sampled within the solution domain  $\Omega$ , given that

$$\|\nabla \psi_p^e\|_2 = 1 \tag{27}$$

### A PiNN with OpenFOAM and PyTorch C++ API

Problem definition II





- A unit-box domain  $x, y, z \in [0, 1]$ , discretized with  $N_c$  cells along each coordinate direction.
- Sphere of radius R = 0.25 centered at  $\mathbf{c} = (0.5, 0.5, 0.5)$ .
- **Motivation**: solving  $\|\nabla \psi_p^e\|_2 = 1$  is quite complex using the unstructured Finite Volume method, if an approximative solution is good-enough, PiNNs are an option.

Implementation: include order in OpenFOAM



```
// libtorch
#include <torch/torch.h>
#include "ATen/Functions.h"
#include "ATen/core/interned_strings.h"
#include "torch/nn/modules/activation.h"
#include "torch/optim/lbfgs.h"
#include "torch/optim/rmsprop.h"
...
// OpenFOAM
#include "fvCFD.H"
```

- Example implemented as an OpenFOAM application that links with the PyTorch C++ API (libtorch).
- Include libtorch before OpenFOAM headers.
- Compile OpenFOAM with C++14 or higher, in wmake/rules/General/C++, set CC = g++\$(COMPILER\_VERSION) -std=c++2a.

Implementation: torch::Sequential Multi-Layer Perceptron



- torch::Sequential used to construct the MLP.
- MLP architecture defined by application options or a dictionary.
  - Parametrization for hyperparameter tuning.

Implementation: OpenFOAM-libtorch GeometricField transfer



```
// - Reinterpret OpenFOAM's input volScalarField as scalar* array
volScalarField::pointer vf_data = vf.ref().data();
// - Use the scalar* (volScalarField::pointer) to view
// the volScalarField as torch::Tensor without copying data.
torch::Tensor vf_tensor = torch::from_blob(vf_data, {vf.size(), 1});
```

- torch::from\_blob constructs a view (interpretation) of data as torch::tensor.
- This works in OpenFOAM because all tensor fields in OpenFOAM are UList<T>

```
template<class T>
class UList
{
    // Private Data
        //- Number of elements in UList
    label size_;
    //- Vector of values of type T
    T* __restrict__ v_;
```

Implementation: PiNN gradient



- Following eq. (18), the scalar-valued forward-pass of the NN is differentiated w.r.t. input data a "tensor" of input points  $\mathbf{x}_p$ .
- For each point  $\mathbf{x}_p = (x, y, z)$  we multiply the Jacobian with 1.
- For efficiency reasons (vectorization), calculation is done for all  $p \in P$ , so we need  $\mathbf{v} = [1, 1, 1, 1, \dots, 1]^T$  in  $J_u \cdot \mathbf{v}$ , the length of  $\mathbf{v}$  is  $N_p$ , the number of points.

Implementation: PiNN loss for our problem



```
// Compute the data mse loss.
auto mse_data = mse_loss(vf_predict, vf_training);

// Compute the gradient mse loss.
auto mse_grad = mse_loss(
    at::norm(vf_predict_grad[0], 2, -1),
    torch::ones_like(vf_training)
);

// Combine the losses into a Physics Informed Neural Network.
mse = mse_data + mse_grad;
```

■ The  $e_{MSF}^r$  is

$$e_{MSE}^{r} = \frac{1}{N_p} \sum_{r=1}^{N_r} (\|\nabla \psi_{\theta}^{nn}(\mathbf{x}_p)\|_2 - 1)^2$$
 (28)

Implementation: hyperparameter tuning I



Grid Search for OpenFOAM+libtorch: works with other methods as well

- A Jupyter Notebook uses **subprocess.call** to run the variations.
- An OpenFOAM application parses hyperparameters as options.
- Results stored as CSV, with repeated hyperparameters as columns

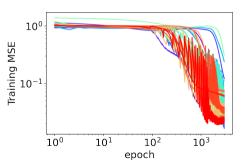
	HIDDEN_LAYERS	OPTIMIZER_STEP	MAX_ITERATIONS	DELTA_X	EPOCH	DATA_MSE	GRAD_MSE	TRAININ
0	10,10,10,10	0.0001	3000	0.0625	1	0.164133	0.927430	1.
1	10,10,10,10	0.0001	3000	0.0625	2	0.154731	0.928236	1.
2	10,10,10,10	0.0001	3000	0.0625	3	0.148419	0.928779	1.
3	10,10,10,10	0.0001	3000	0.0625	4	0.143443	0.929211	1.

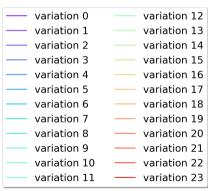
- CSVs are concatenated and the final **pandas.DataFrame** can easily be filtered by **repeated** hyperparameters. Avoid the **pandas.MultiIndex**, it's not worth the trouble.
- Each parameter vector is uniquely identified with the ID that is part of the file name: pinnFoam-00000000.csv - connected to the hyperparameters using pandas.DataFrame.unique on hyperparameter columns.

Implementation: hyperparameter tuning II



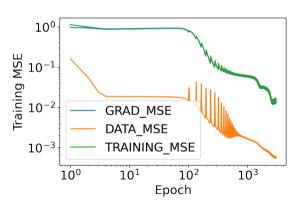
#### Loss MSE Grid Search





Implementation: best hyperparameters





#### Best hyperparameters:

HIDDEN\_LAYERS OPTIMIZER\_STEP 20,20,20,20 0.001 TRAINING\_ID 3000 3

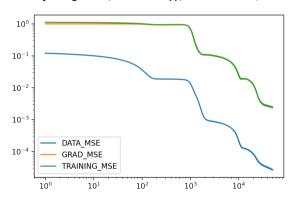
- Just adding different  $e_{\it MSE}$  leads to different gradient flows.
- These are example runs, actual runs have way more epochs.
- **Exercise**: try to find  $\lambda_r$  and  $\lambda_d$  in

$$l_{MSE} = \lambda_d e^d_{MSE} + \lambda_r e^r_{MSE}$$

Implementation: real-world training

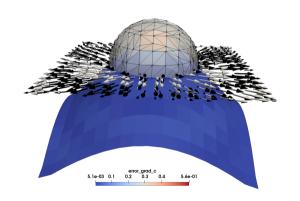


Real-world training takes a very long time, minutes (!), even with  $N_c = 16$  (4096 cells!).



Implementation: visualization





- Wireframe:  $\psi_{\theta^M}^{nn}$  iso-surface.
- Gray surface:  $\psi^e$  iso-surface.
- Black and gray arrows:  $(\nabla \psi)^e$  and  $\nabla \psi_{\theta^M}^{nn}$  gradient vectors.
- Warped surface:  $\|(\nabla \psi)^e \nabla \psi_{\theta^M}^{nn}\|_2$

Source Code and Data



- https://gitlab.com/tmaric/ofw17-training-physics-based-dl
- Source Code and Data snapshot [Maric, 2022-07-10].

## (My personal) Conclusions on PiNNs



- A very promising tool for some challenging problems.
- PiNNs aren't replacing standard numerics any time soon for standard problems.
- Very simple compared to standard numerics: no need to implement very complex algorithms for solving complex (systems of) PDEs.
- Know your Jacobians for PDE differential operator calculation using Automatic Differentiation.
- The PDE differential operators require differentiable activation functions difficult to train.
- If something goes wrong, guess again numerical methods can use more analysis.
- Hyperparameter tuning takes a lot of effort: automation makes sense.

#### References I



- Andreas Griewank and Andrea Walther. Evaluating derivatives: principles and techniques of algorithmic differentiation. SIAM, 2008.
- Ameya D. Jagtap, Kenji Kawaguchi, and George Em Karniadakis. Adaptive activation functions accelerate convergence in deep and physics-informed neural networks. *Journal of Computational Physics*, 404:109136, 2020. ISSN 10902716. doi: 10.1016/j.jcp.2019.109136. URL https://doi.org/10.1016/j.jcp.2019.109136. arXiv: 1906.01170 Publisher: Elsevier Inc.
- I.E. Lagaris, A.C. Likas, and D.G. Papageorgiou. Neural-network methods for boundary value problems with irregular boundaries. *IEEE Transactions on Neural Networks*, 11(5):1041–1049, September 2000. ISSN 10459227. doi: 10.1109/72.870037. URL http://ieeexplore.ieee.org/document/870037/.
- Isaac Elias Lagaris, Aristidis Likas, and Dimitrios I Fotiadis. Artificial Neural Networks for Solving Ordinary and Partial Differential Equations. *IEEE Transactions on Neural Networks*, 9(5):987–1000, 1998. doi: 10.1109/72.712178.
- Tomislav Maric. Towards physics-based deep learning in openfoam: Combining openfoam with the pytorch c++ api (source code and data), 2022-07-10. URL https://tudatalib.ulb.tu-darmstadt.de/handle/tudatalib/3527.

#### References II



K.S. McFall and J.R. Mahan. Artificial Neural Network Method for Solution of Boundary Value Problems With Exact Satisfaction of Arbitrary Boundary Conditions. *IEEE Transactions on Neural Networks*, 20(8):1221–1233, August 2009. ISSN 1045-9227, 1941-0093. doi: 10.1109/TNN.2009.2020735. URL http://ieeexplore.ieee.org/document/5061501/.

Michael L Overton. Numerical computing with IEEE floating point arithmetic. SIAM, 2001.

Apostolos F Psaros, Kenji Kawaguchi, and George Em Karniadakis. Meta-learning PINN loss functions. *Journal of Computational Physics*, 458:111121, June 2022. ISSN 00219991. doi: 10.1016/j.jcp.2022.111121. URL https://linkinghub.elsevier.com/retrieve/pii/S0021999122001838.

M. Raissi, P. Perdikaris, and G. E. Karniadakis. Physics-informed neural networks: A deep learning framework for solving forward and inverse problems involving nonlinear partial differential equations. *Journal of Computational Physics*, 378: 686–707, 2019. ISSN 10902716. doi: 10.1016/j.jcp.2018.10.045. URL https://doi.org/10.1016/j.jcp.2018.10.045. Publisher: Elsevier Inc.

Justin Sirignano and Konstantinos Spiliopoulos. DGM: A deep learning algorithm for solving partial differential equations. Journal of Computational Physics, 375:1339-1364, December 2018. ISSN 00219991. doi: 10.1016/j.jcp.2018.08.029. URL https://linkinghub.elsevier.com/retrieve/pii/S0021999118305527.

#### References III



N Thuerey, P Holl, M Mueller, P Schnell, F Trost, and K Um. Physics-based Deep Learning. 2022. arXiv: 2109.05237v2. Sifan Wang, Yujun Teng, and Paris Perdikaris. Understanding and Mitigating Gradient Flow Pathologies in Physics-Informed Neural Networks. SIAM Journal on Scientific Computing, 43(5):A3055-A3081, January 2021. ISSN 1064-8275. doi: 10.1137/20M1318043. URL https://epubs.siam.org/doi/10.1137/20M1318043. Publisher: Society for Industrial and Applied Mathematics.